

Singularities, Shocks in Hele-Shaw flow

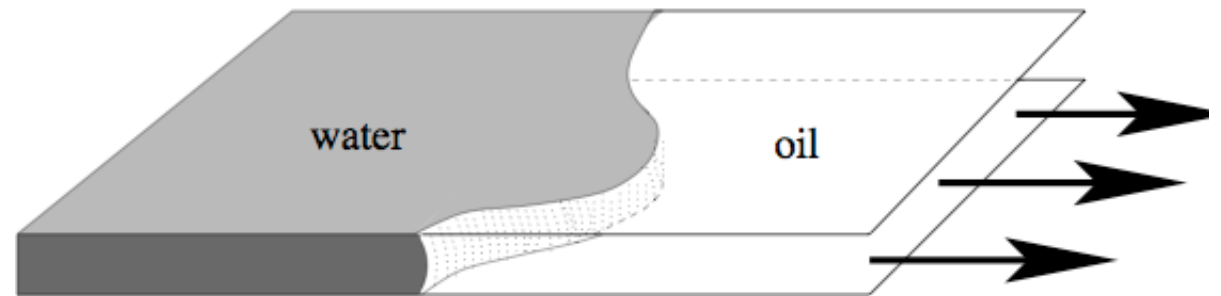
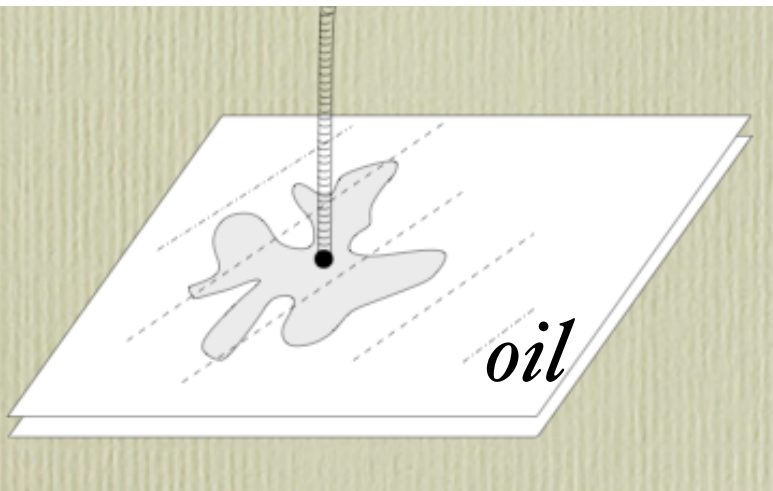
as

Stokes phenomena in Painlevé I

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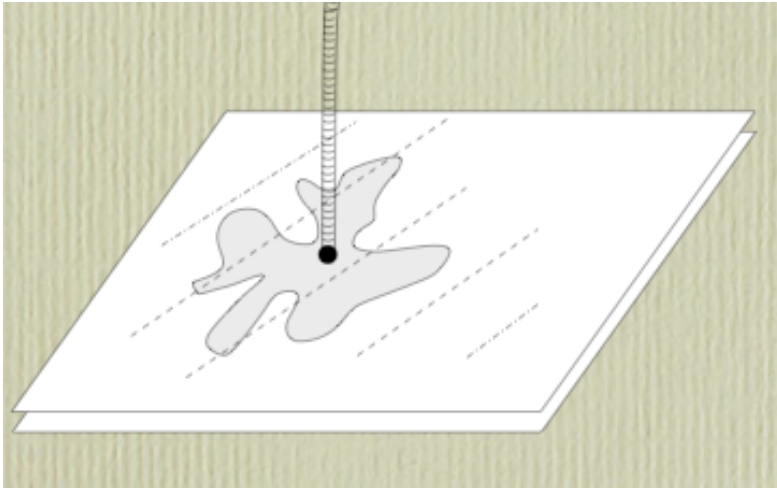
Hele-Shaw cell (1894)

water



Oil (*exterior*)-incompressible liquid with high viscosity

Water (*interior*) - incompressible liquid with low viscosity



$$\mathbf{v} = -\nabla p$$

$$\Delta p = 0$$

$$p|_{\Gamma} = 0$$

$$p|_{\infty} = -\log |z|$$

Darcy Law

velocity of a moving planar interface =
= a gradient of a harmonic field

Painlevé I

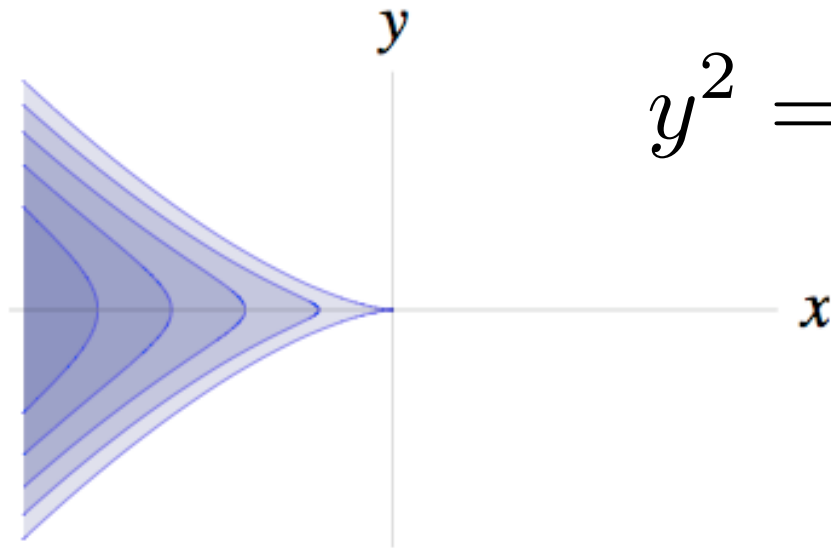
$$u_{tt} - 6u^2 = \hbar t.$$

Physical Solution: $u \rightarrow +\sqrt{-\hbar t/6}$

$$\Psi(X, t) \quad \left(\partial_t^2 - 2u(t) \right) \Psi = X \Psi$$

$$X = x + iy \quad T = \frac{\hbar t}{6}$$

Finite time cusp singularities



$$y^2 = -(x + 2u(t)) \left(x - u(t) \right)^2$$

$$u(t) = \sqrt{-t}$$

$$y^2 \sim x^3$$

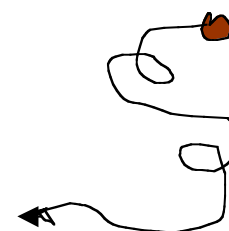
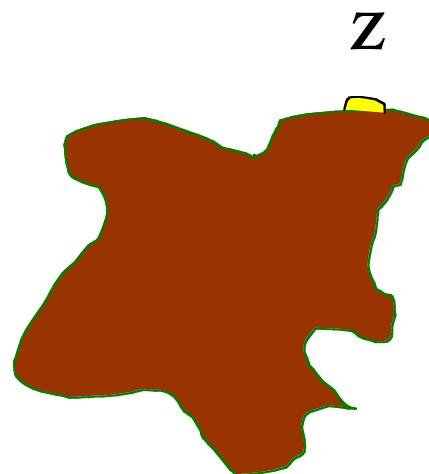
Brownian excursion with a moving boundary

A probability for BM to arrive on an element of the boundary is a *harmonic measure* of the boundary

$$\Delta p = 0, \quad p(z \rightarrow \infty) \rightarrow \log |z|$$

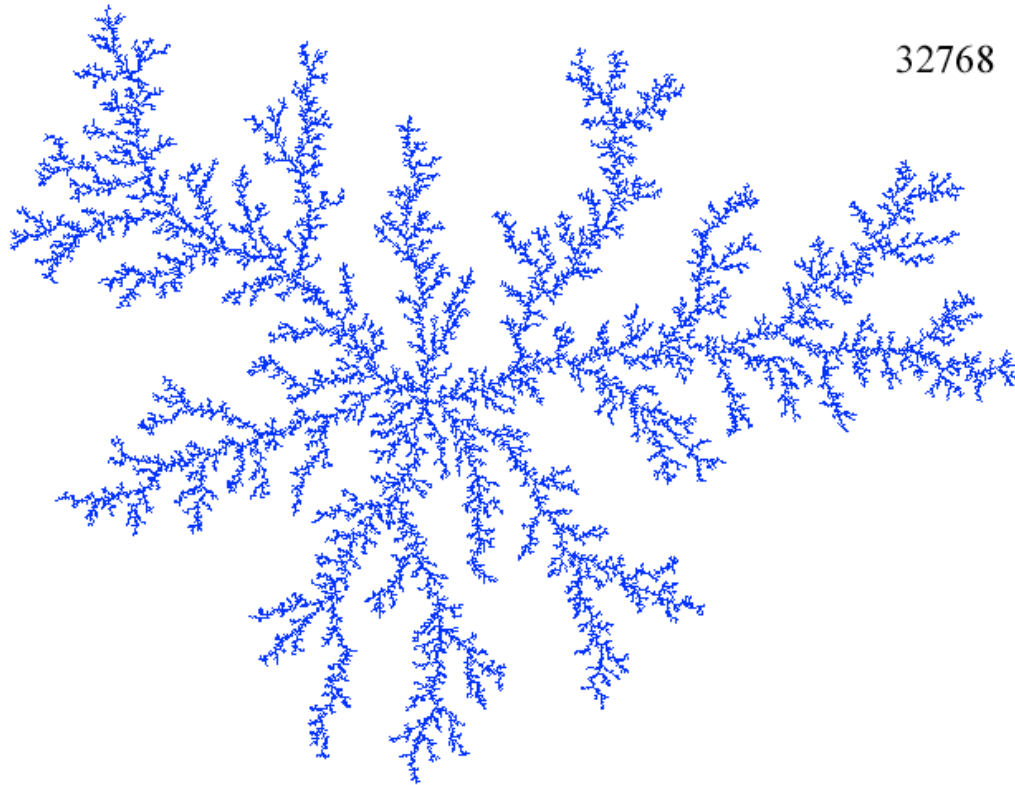
$|\nabla p|$ – harmonic measure

$p = 0$ – on the boundary



Continuous problem (a hydrodynamic limit):

a size of particles $\hbar \rightarrow 0$ tends to zero



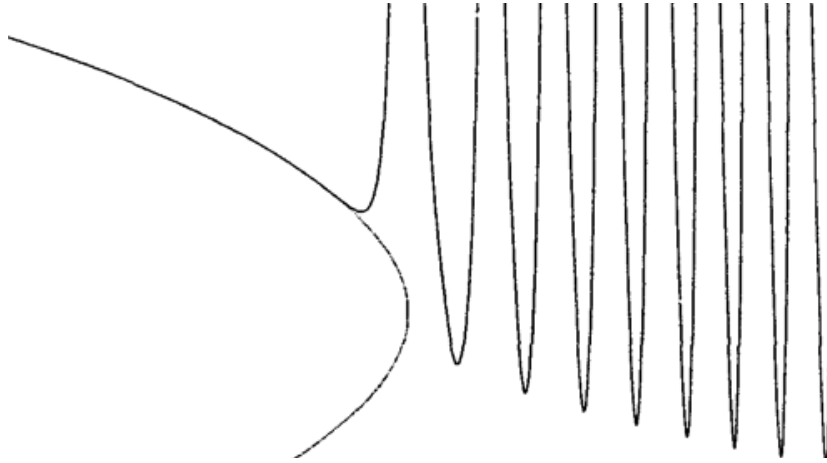
zeros of $\Psi(X, t)$

The limit : size of particles tends to zero is ill-defined

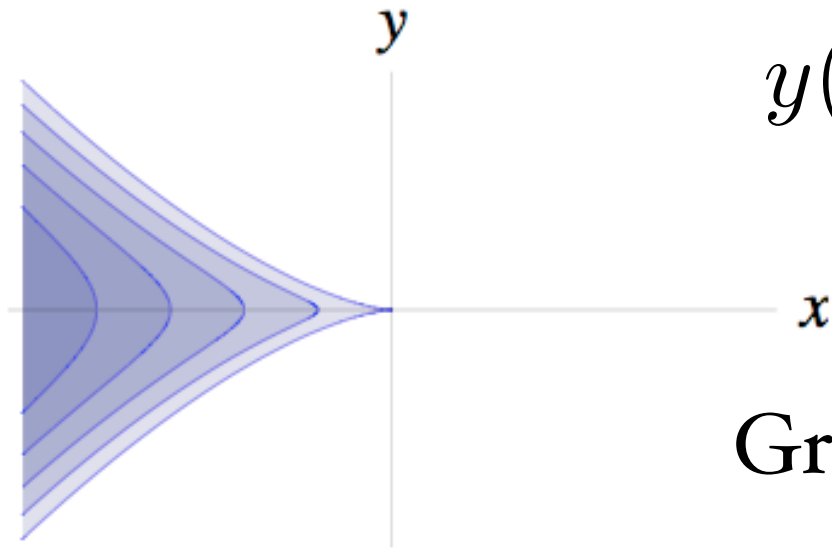
Weak solution

Darcy law is applied everywhere in the fluid except on a moving, growing and branching graph of shocks, where pressure suffers a finite discontinuity.

$$u_{tt} - 6u^2 = \hbar t.$$



Evolution of an elliptic curve



$$y(x) \rightarrow Y(X)$$

Graph to a Complex Curve

Darcy Law has a form of
Witham equation:

Zabrodin, Krichever, P.W.

$$\partial_T Y = -\partial_X \phi,$$

$$\phi(X) = \psi + ip$$

Integral Form of the Darcy Law

$$\begin{aligned}\frac{d}{dT} \operatorname{Im} \oint_B (-iY) dX &= \oint_B \mathbf{v} \times \mathbf{d}\ell = \oint_B d\psi = \text{flux of fluid}, \\ \frac{d}{dT} \operatorname{Re} \oint_B (-iY) dX &= - \oint_B \mathbf{v} \cdot \mathbf{d}\ell = - \oint_B dp = \text{vorticity}.\end{aligned}$$

$$\operatorname{Re} \oint_{all\ cycles} (-iY) dX = 0$$

$$\operatorname{Im} \oint_{\infty} (-iY) dX = \text{time}$$

A unique weak solution

$$\operatorname{Re} \oint_{all \text{ cycles}} (-iY) dX = 0$$

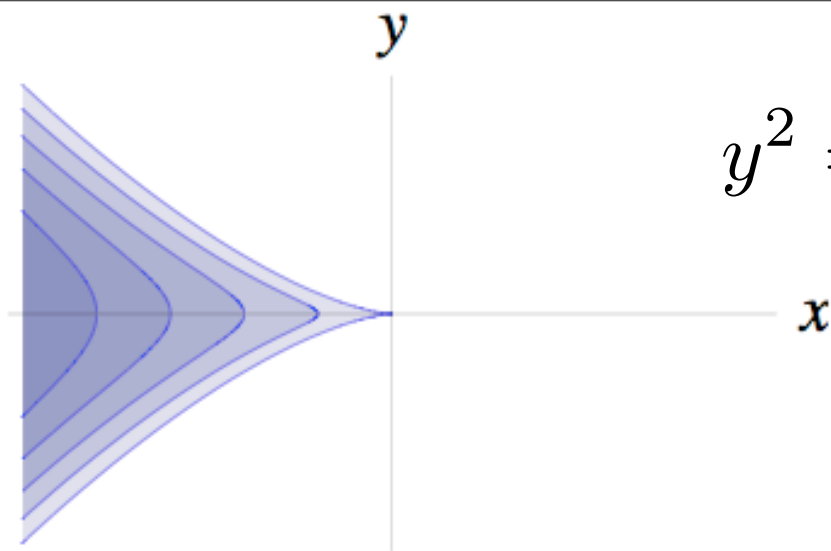
$$\operatorname{Im} \oint_{\infty} (-iY) dX = \text{time}$$

Boutroux condition/ Krichever

Shocks and anti-Stokes lines

$$\operatorname{Re} \int^X (-iY) dX = 0, \quad X \in \text{shocks}$$

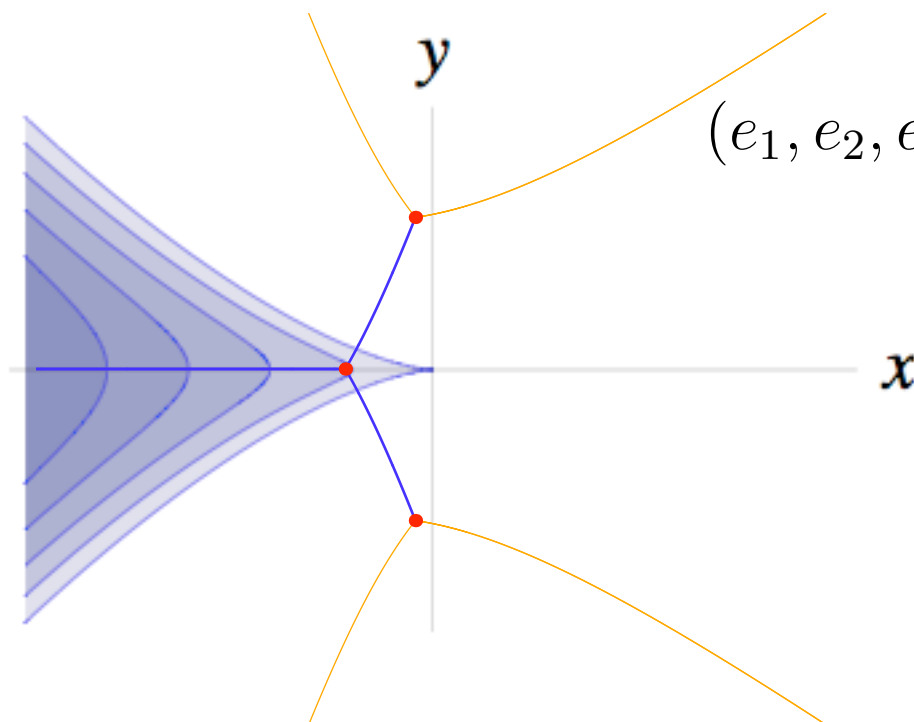
$$\operatorname{Re} \int^X (-iY) dX > 0, \quad X \rightarrow \text{shocks}$$



$$y^2 = -(x + 2u(t))(x - u(t))^2$$

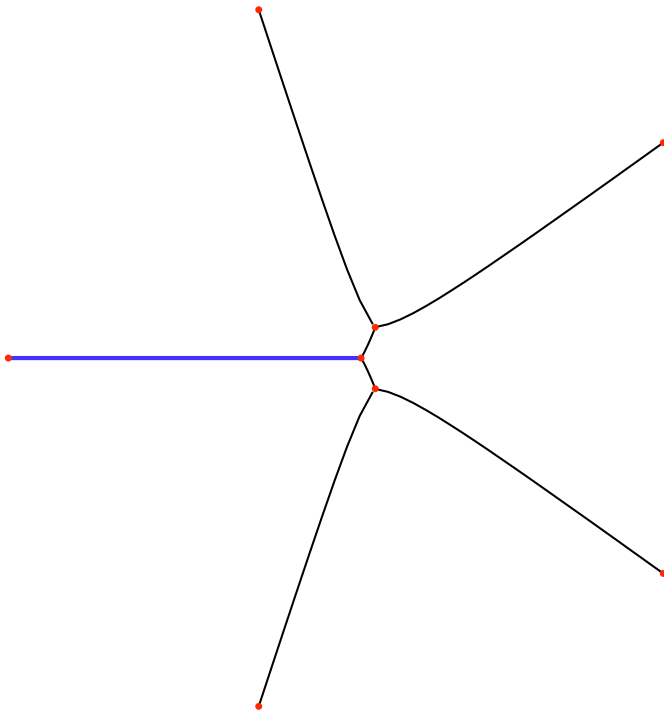
$$u(t) = \sqrt{-t}$$

$$y^2 = (x - e_1)(x - e_2)(x - e_3)$$

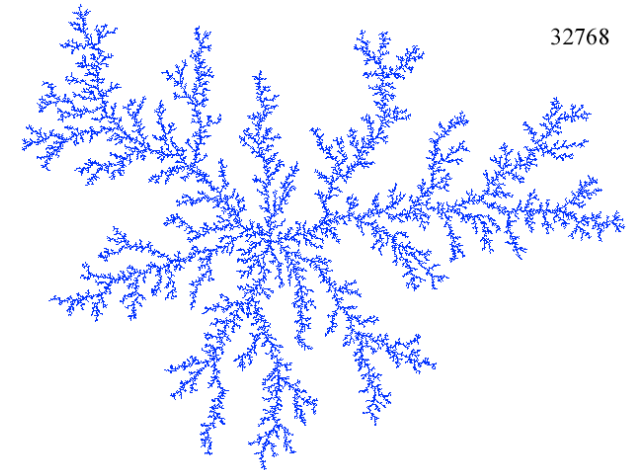


$$(e_1, e_2, e_3) = \sqrt{\frac{3}{h^2 - \frac{3}{4}}}(-1, \frac{1}{2} + ih, \frac{1}{2} - ih)\sqrt{t}$$

$$h \approx 3.246382253744278875676.$$



Two generations of branching



32768 generation of branching

Stokes Phenomenon in Painlevé I

Painlevé I

$$u_{tt} - 6u^2 = \hbar t.$$

$$\left(\partial_t^2 - 2u(t)\right)\Psi = X\Psi$$

$$Y(X, t) = \lim_{\hbar \rightarrow 0} (i\hbar \partial_X) \log \Psi(X, t)$$

$$\hbar \log \Psi|_{X \rightarrow \infty} \rightarrow \frac{4}{5} X^{5/2} - t X^{1/2}, \quad |\text{Arg}|X| < \pi.$$

$$\psi_{\pm}^{WKB}(X, t) \sim e^{\pm \frac{1}{\hbar} \int^X (-iY) dX}$$

Time-independent monodromy data:

$$\operatorname{Re} \oint_{all \text{ cycles}} (-iY) dX = 0$$

$$\operatorname{Im} \oint_{\infty} (-iY) dX = \text{time}$$

Boutroux, Krichever, Novikov-Grinevich

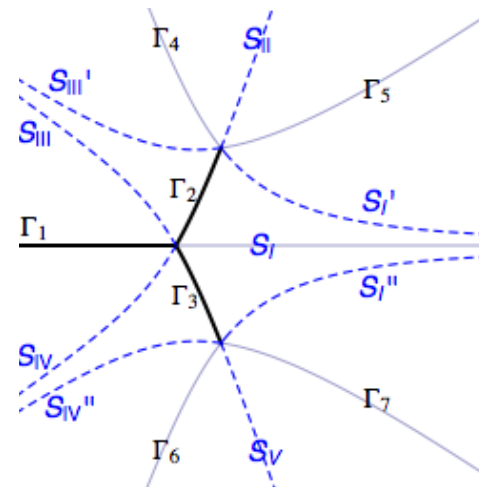
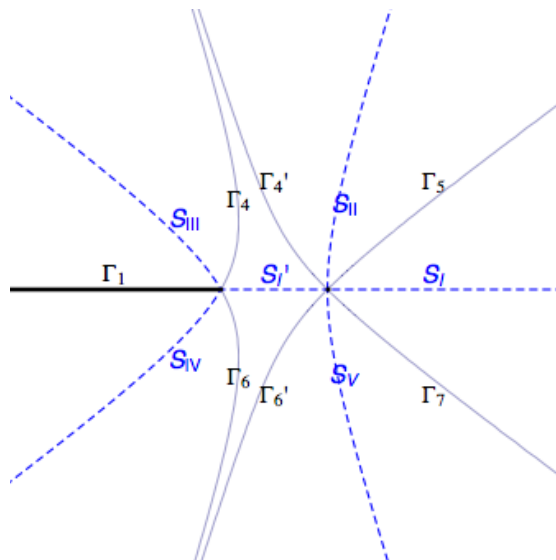
identical to the Darcy law

Darcy law requires the asymptote

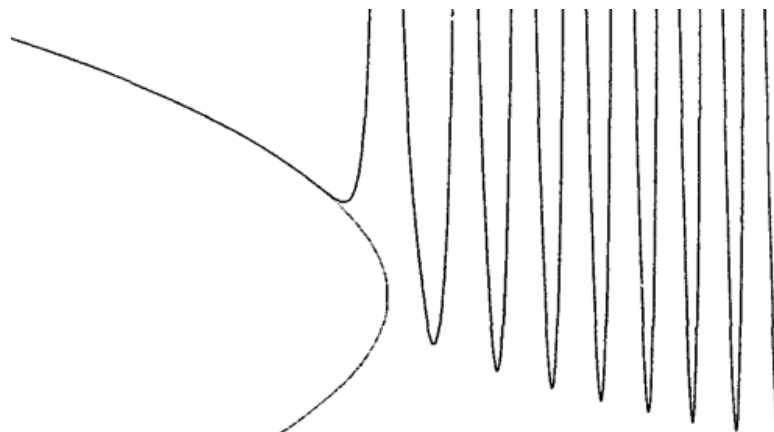
$$\hbar \log \Psi|_{X \rightarrow \infty} \rightarrow \frac{4}{5} X^{5/2} - t X^{1/2}, \quad |\text{Arg}|X| < \pi.$$

$$\Psi^{WKB} = A_S \psi_+^{WKB} + B_S \psi_-^{WKB},$$

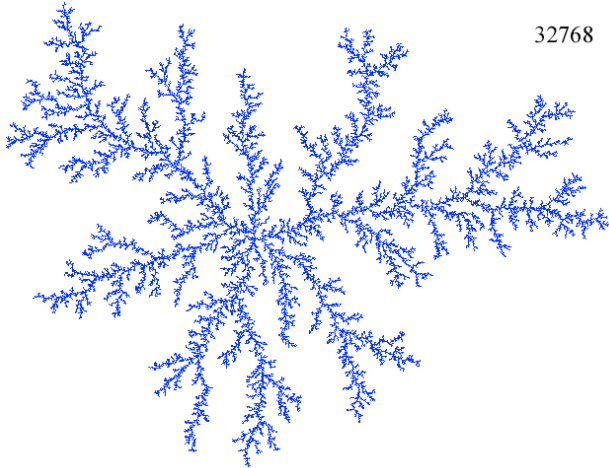
$$\psi_{\pm}^{WKB}(X, t) \sim e^{\pm \frac{1}{\hbar} \int^X (-iY) dX}$$



Position of zeros of the physical wave-function



Smooth and oscillatory regimes of the Painleve function



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Hypothesis:

The pattern is Stokes lines of higher order Painleve equations.